

STAT 801: Mathematical Statistics

Inversion of Generating Functions

Previous theorem is non-constructive characterization. Can get from ϕ_X to F_X or f_X by **inversion**. See homework for basic **inversion** formula:

If X is a random variable taking only integer values then for each integer k

$$\begin{aligned} P(X = k) &= \frac{1}{2\pi} \int_0^{2\pi} \phi_X(t) e^{-itk} dt \\ &= \frac{1}{2\pi} \int_{-\pi}^{\pi} \phi_X(t) e^{-itk} dt. \end{aligned}$$

The proof proceeds from the formula

$$\phi_X(t) = \sum_k e^{ikt} P(X = k).$$

Now suppose that X has a continuous bounded density f . Define

$$X_n = [nX]/n$$

where $[a]$ denotes the integer part (rounding down to the next smallest integer). We have

$$\begin{aligned} P(k/n \leq X < (k+1)/n) &= P([nX] = k) \\ &= \frac{1}{2\pi} \int_{-\pi}^{\pi} \phi_{[nX]}(t) \\ &\quad \times e^{-itk} dt. \end{aligned}$$

Make the substitution $t = u/n$, and get

$$nP(k/n \leq X < (k+1)/n) = \frac{1}{2\pi} \times \int_{-n\pi}^{n\pi} \phi_{[nX]}(u/n) e^{iuk/n} du$$

Now, as $n \rightarrow \infty$ we have

$$\phi_{[nX]}(u/n) = E(e^{iu[nX]/n}) \rightarrow E(e^{iuX})$$

(by the dominated convergence theorem – the dominating random variable is just the constant 1). The range of integration converges to the whole real line and if $k/n \rightarrow x$ we see that the left hand side converges to the density $f(x)$ while the right hand side converges to

$$\frac{1}{2\pi} \int_{-\infty}^{\infty} \phi_X(u) e^{-iux} du$$

which gives the inversion formula

$$f_X(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \phi_X(u) e^{-iux} du$$

Many other such formulas are available to compute things like $F(b) - F(a)$ and so on.

All such formulas are sometimes referred to as Fourier inversion formulas; the characteristic function itself is sometimes called the Fourier transform of the distribution or cdf or density of X .

Inversion of the Moment Generating Function

MGF and characteristic function related formally:

$$M_X(it) = \phi_X(t)$$

When M_X exists this relationship is not merely formal; the methods of complex variables mean there is a “nice” (analytic) function which is $E(e^{zX})$ for any complex $z = x + iy$ for which $M_X(x)$ is finite.

SO: there is an inversion formula for M_X using a complex *contour integral*:

If z_1 and z_2 are two points in the complex plane and C a path between these two points we can define the path integral

$$\int_C f(z)dz$$

by the methods of line integration.

Do algebra with such integrals via usual theorems of calculus. The Fourier inversion formula was

$$2\pi f(x) = \int_{-\infty}^{\infty} \phi(t)e^{-itx} dt$$

so replacing ϕ by M we get

$$2\pi f(x) = \int_{-\infty}^{\infty} M(it)e^{-itx} dt$$

If we just substitute $z = it$ then we find

$$2\pi i f(x) = \int_C M(z)e^{-zx} dz$$

where the path C is the imaginary axis. Methods of complex integration permit us to replace C by any other path which starts and ends at the same place. Sometimes can choose path to make it easy to do the integral approximately; this is what saddlepoint approximations are. Inversion formula is called the inverse Laplace transform; the mgf is also called the Laplace transform of the distribution or cdf or density.

Applications of Inversion

1): Numerical calculations

Example: Many statistics have a distribution which is approximately that of

$$T = \sum \lambda_j Z_j^2$$

where the Z_j are iid $N(0, 1)$. In this case

$$\begin{aligned} E(e^{itT}) &= \prod E(e^{it\lambda_j Z_j^2}) \\ &= \prod (1 - 2it\lambda_j)^{-1/2}. \end{aligned}$$

Imhof (*Biometrika*, 1961) gives a simplification of the Fourier inversion formula for

$$F_T(x) - F_T(0)$$

which can be evaluated numerically:

$$\begin{aligned} F_T(x) - F_T(0) &= \int_0^x f_T(y) dy \\ &= \int_0^x \frac{1}{2\pi} \int_{-\infty}^{\infty} \prod (1 - 2it\lambda_j)^{-1/2} e^{-ity} dt dy \end{aligned}$$

Multiply

$$\phi(t) = \left[\frac{1}{\prod(1 - 2it\lambda_j)} \right]^{1/2}$$

top and bottom by the complex conjugate of the denominator:

$$\phi(t) = \left[\frac{\prod(1 + 2it\lambda_j)}{\prod(1 + 4t^2\lambda_j^2)} \right]^{1/2}$$

The complex number $1 + 2it\lambda_j$ is $r_j e^{i\theta_j}$ where $r_j = \sqrt{1 + 4t^2\lambda_j^2}$ and $\tan(\theta_j) = 2t\lambda_j$. This allows us to rewrite

$$\phi(t) = \left[\frac{\prod r_j e^{i\sum\theta_j}}{\prod r_j^2} \right]^{1/2}$$

or

$$\phi(t) = \frac{e^{i\sum \tan^{-1}(2t\lambda_j)/2}}{\prod(1 + 4t^2\lambda_j^2)^{1/4}}$$

Assemble this to give

$$F_T(x) - F_T(0) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{e^{i\theta(t)}}{\rho(t)} \int_0^x e^{-iyt} dy dt$$

where

$$\theta(t) = \sum \tan^{-1}(2t\lambda_j)/2$$

and $\rho(t) = \prod(1 + 4t^2\lambda_j^2)^{1/4}$. But

$$\int_0^x e^{-iyt} dy = \frac{e^{-ixt} - 1}{-it}$$

We can now collect up the real part of the resulting integral to derive the formula given by Imhof. I don't produce the details here.

2): The central limit theorem (in some versions) can be deduced from the Fourier inversion formula: if X_1, \dots, X_n are iid with mean 0 and variance 1 and $T = n^{1/2}\bar{X}$ then with ϕ denoting the characteristic function of a single X we have

$$\begin{aligned} E(e^{itT}) &= E(e^{in^{-1/2}t\sum X_j}) \\ &= [\phi(n^{-1/2}t)]^n \\ &\approx \left[\phi(0) + \frac{t\phi'(0)}{\sqrt{n}} + \frac{t^2\phi''(0)}{2n} + o(n^{-1}) \right]^n \end{aligned}$$

But now $\phi(0) = 1$ and

$$\phi'(t) = \frac{d}{dt} E(e^{itX_1}) = iE(X_1 e^{itX_1})$$

So $\phi'(0) = E(X_1) = 0$. Similarly

$$\phi''(t) = i^2 E(X_1^2 e^{itX_1})$$

so that

$$\phi''(0) = -E(X_1^2) = -1$$

It now follows that

$$E(e^{itT}) \approx [1 - t^2/(2n) + o(1/n)]^n \rightarrow e^{-t^2/2}.$$

With care we can then apply the Fourier inversion formula and get

$$\begin{aligned} f_T(x) &= \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{-itx} [\phi(tn^{-1/2})]^n dt \\ &\rightarrow \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{-itx} e^{-t^2/2} dt \\ &= \frac{1}{\sqrt{2\pi}} \phi_Z(-x) \end{aligned}$$

where ϕ_Z is the characteristic function of a standard normal variable Z . Doing the integral we find

$$\phi_Z(x) = \phi_Z(-x) = e^{-x^2/2}$$

so that

$$f_T(x) \rightarrow \frac{1}{\sqrt{2\pi}} e^{-x^2/2}$$

which is a standard normal random variable.

This proof of the central limit theorem is not terribly general since it requires T to have a bounded continuous density. The central limit theorem itself is a statement about cdfs not densities and is

$$P(T \leq t) \rightarrow P(Z \leq t).$$

3) Saddlepoint approximation from MGF inversion formula

$$2\pi i f(x) = \int_{-i\infty}^{i\infty} M(z) e^{-zx} dz$$

(limits of integration indicate contour integral running up imaginary axis.) Replace contour (using complex variables) with line $Re(z) = c$. ($Re(Z)$ denotes the real part of z , that is, x when $z = x + iy$ with x and y real.) Must choose c so that $M(c) < \infty$. Rewrite inversion formula using cumulant generating function $K(t) = \log(M(t))$:

$$2\pi i f(x) = \int_{c-i\infty}^{c+i\infty} \exp(K(z) - zx) dz.$$

Along the contour in question we have $z = c + iy$ so we can think of the integral as being

$$i \int_{-\infty}^{\infty} \exp(K(c + iy) - (c + iy)x) dy$$

Now do a Taylor expansion of the exponent:

$$K(c + iy) - (c + iy)x = K(c) - cx + iy(K'(c) - x) - y^2 K''(c)/2 + \dots$$

Ignore the higher order terms and select a c so that the first derivative

$$K'(c) - x$$

vanishes. Such a c is a saddlepoint. We get the formula

$$2\pi f(x) \approx \exp(K(c) - cx) \times \int_{-\infty}^{\infty} \exp(-y^2 K''(c)/2) dy.$$

The integral is just a normal density calculation and gives $\sqrt{2\pi/K''(c)}$. The saddlepoint approximation is

$$f(x) = \frac{\exp(K(c) - cx)}{\sqrt{2\pi K''(c)}}.$$

Essentially the same idea lies at the heart of the proof of Sterling's approximation to the factorial function:

$$n! = \int_0^{\infty} \exp(n \log(x) - x) dx$$

The exponent is maximized when $x = n$. For n large we approximate $f(x) = n \log(x) - x$ by

$$f(x) \approx f(x_0) + (x - x_0)f'(x_0) + (x - x_0)^2 f''(x_0)/2$$

and choose $x_0 = n$ to make $f'(x_0) = 0$. Then

$$n! \approx \int_0^{\infty} \exp[n \log(n) - n - (x - n)^2/(2n)] dx$$

Substitute $y = (x - n)/\sqrt{n}$ to get the approximation

$$n! \approx n^{1/2} n^n e^{-n} \int_{-\infty}^{\infty} e^{-y^2/2} dy$$

or

$$n! \approx \sqrt{2\pi n} n^{n+1/2} e^{-n}$$

This tactic is called Laplace's method. Note that I am being very sloppy about the limits of integration; to do the thing properly you have to prove that the integral over x not near n is negligible.