

## Inversion

Previous theorem is non-constructive characterization.

Can get from  $\phi_X$  to  $F_X$  or  $f_X$  by **inversion**.

See homework for basic **inversion** formula:

If  $X$  is a random variable taking only integer values then for each integer  $k$

$$\begin{aligned} P(X = k) &= \frac{1}{2\pi} \int_0^{2\pi} \phi_X(t) e^{-itk} dt \\ &= \frac{1}{2\pi} \int_{-\pi}^{\pi} \phi_X(t) e^{-itk} dt. \end{aligned}$$

The proof proceeds from the formula

$$\phi_X(t) = \sum_k e^{ikt} P(X = k).$$

Now suppose  $X$  has continuous bounded density  $f$ . Define

$$X_n = [nX]/n$$

where  $[a]$  denotes the integer part (rounding down to the next smallest integer). We have

$$\begin{aligned} P(k/n \leq X < (k+1)/n) \\ &= P([nX] = k) \\ &= \frac{1}{2\pi} \int_{-\pi}^{\pi} \phi_{[nX]}(t) \times e^{-itk} dt. \end{aligned}$$

Make the substitution  $t = u/n$ , and get

$$\begin{aligned} nP(k/n \leq X < (k+1)/n) &= \frac{1}{2\pi} \\ &\times \int_{-n\pi}^{n\pi} \phi_{[nX]}(u/n) e^{-iuk/n} du. \end{aligned}$$

Now, as  $n \rightarrow \infty$  we have

$$\phi_{[nX]}(u/n) = E(e^{iu[nX]/n}) \rightarrow E(e^{iuX}).$$

(Dominated convergence:  $|e^{iu}| \leq 1$ .)

Range of integration converges to the whole real line.

If  $k/n \rightarrow x$  left hand side converges to density  $f(x)$  while right hand side converges to

$$\frac{1}{2\pi} \int_{-\infty}^{\infty} \phi_X(u) e^{-iux} du$$

which gives the inversion formula

$$f_X(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \phi_X(u) e^{-iux} du.$$

Many other such formulas are available to compute things like  $F(b) - F(a)$  and so on.

All such formulas called **Fourier inversion formulas**.

Characteristic ftn also called **Fourier transform** of  $f$  or  $F$ .

## Inversion of the Moment Generating Function

MGF and characteristic function related formally:

$$M_X(it) = \phi_X(t).$$

When  $M_X$  exists this relationship is not merely formal; the methods of complex variables mean there is a “nice” (analytic) function which is  $E(e^{zX})$  for any complex  $z = x + iy$  for which  $M_X(x)$  is finite.

SO: there is an inversion formula for  $M_X$  using a complex *contour integral*:

If  $z_1$  and  $z_2$  are two points in the complex plane and  $C$  a path between these two points we can define the path integral

$$\int_C f(z) dz$$

by the methods of line integration.

Do algebra with such integrals via usual theorems of calculus.

The Fourier inversion formula was

$$2\pi f(x) = \int_{-\infty}^{\infty} \phi(t)e^{-itx} dt$$

so replacing  $\phi$  by  $M$  we get

$$2\pi f(x) = \int_{-\infty}^{\infty} M(it)e^{-itx} dt.$$

If we just substitute  $z = it$  then we find

$$2\pi i f(x) = \int_C M(z)e^{-zx} dz$$

where the path  $C$  is the imaginary axis.

Complex contour integration: replace  $C$  by any other path which starts and ends at the same place.

Sometimes can choose path to make it easy to do the integral approximately; this is what **saddlepoint approximations** are.

Inversion formula is called the **inverse Laplace transform**; the mgf is also called the Laplace transform of  $f$  or  $F$ .

## Applications of Inversion

1): Numerical calculations

Example: Many statistics have a distribution which is approximately that of

$$T = \sum \lambda_j Z_j^2$$

where the  $Z_j$  are iid  $N(0, 1)$ . In this case

$$\begin{aligned} E(e^{itT}) &= \prod E(e^{it\lambda_j Z_j^2}) \\ &= \prod (1 - 2it\lambda_j)^{-1/2}. \end{aligned}$$

Imhof (*Biometrika*, 1961) gives a simplification of the Fourier inversion formula for

$$F_T(x) - F_T(0)$$

which can be evaluated numerically:

$$\begin{aligned} F_T(x) - F_T(0) &= \int_0^x f_T(y) dy \\ &= \int_0^x \frac{1}{2\pi} \int_{-\infty}^{\infty} \prod (1 - 2it\lambda_j)^{-1/2} e^{-ity} dt dy. \end{aligned}$$

Multiply

$$\phi(t) = \left[ \frac{1}{\prod(1 - 2it\lambda_j)} \right]^{1/2}$$

top and bottom by the complex conjugate of the denominator:

$$\phi(t) = \left[ \frac{\prod(1 + 2it\lambda_j)}{\prod(1 + 4t^2\lambda_j^2)} \right]^{1/2} .$$

The complex number  $1 + 2it\lambda_j$  is  $r_j e^{i\theta_j}$  where

$$r_j = \sqrt{1 + 4t^2\lambda_j^2}$$

and

$$\tan(\theta_j) = 2t\lambda_j .$$

This allows us to rewrite

$$\phi(t) = \left[ \frac{\prod r_j e^{i \sum \theta_j}}{\prod r_j^2} \right]^{1/2}$$

or

$$\phi(t) = \frac{e^{i \sum \tan^{-1}(2t\lambda_j)/2}}{\prod (1 + 4t^2\lambda_j^2)^{1/4}}.$$

Assemble this to give

$$F_T(x) - F_T(0) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{e^{i\theta(t)}}{\rho(t)} \int_0^x e^{-iyt} dy dt$$

where

$$\theta(t) = \sum \tan^{-1}(2t\lambda_j)/2$$

and  $\rho(t) = \prod (1 + 4t^2\lambda_j^2)^{1/4}$ . But

$$\int_0^x e^{-iyt} dy = \frac{e^{-ixt} - 1}{-it}.$$

We can now collect up the real part of the resulting integral to derive the formula given by Imhof. I don't produce the details here.



2): The central limit theorem (in some versions) can be deduced from the Fourier inversion formula: if  $X_1, \dots, X_n$  are iid with mean 0 and variance 1 and  $T = n^{1/2}\bar{X}$  then with  $\phi$  denoting the characteristic function of a single  $X$  we have

$$\begin{aligned} E(e^{itT}) &= E(e^{in^{-1/2}t \sum X_j}) \\ &= [\phi(n^{-1/2}t)]^n \\ &\approx \left[ \phi(0) + \frac{t\phi'(0)}{\sqrt{n}} + \frac{t^2\phi''(0)}{2n} + o(n^{-1}) \right]^n \end{aligned}$$

But now  $\phi(0) = 1$  and

$$\phi'(t) = \frac{d}{dt}E(e^{itX_1}) = iE(X_1e^{itX_1}).$$

So  $\phi'(0) = E(X_1) = 0$ . Similarly

$$\phi''(t) = i^2E(X_1^2e^{itX_1})$$

so that

$$\phi''(0) = -E(X_1^2) = -1.$$

It now follows that

$$\begin{aligned} E(e^{itT}) &\approx [1 - t^2/(2n) + o(1/n)]^n \\ &\rightarrow e^{-t^2/2}. \end{aligned}$$

With care we can then apply the Fourier inversion formula and get

$$\begin{aligned} f_T(x) &= \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{-itx} [\phi(tn^{-1/2})]^n dt \\ &\rightarrow \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{-itx} e^{-t^2/2} dt \\ &= \frac{1}{\sqrt{2\pi}} \phi_Z(-x) \end{aligned}$$

where  $\phi_Z$  is the characteristic function of a standard normal variable  $Z$ . Doing the integral we find

$$\phi_Z(x) = \phi_Z(-x) = e^{-x^2/2}$$

so that

$$f_T(x) \rightarrow \frac{1}{\sqrt{2\pi}} e^{-x^2/2}$$

which is a standard normal density.

Proof of the central limit theorem not general: requires  $T$  to have bounded continuous density.

Central limit theorem: statement about cdfs not densities:

$$P(T \leq t) \rightarrow P(Z \leq t).$$

**3)** Saddlepoint approximation from MGF inversion formula

$$2\pi i f(x) = \int_{-i\infty}^{i\infty} M(z) e^{-zx} dz$$

(limits of integration indicate contour integral running up imaginary axis.)

Replace contour (using complex variables) with line  $Re(z) = c$ . ( $Re(Z)$  denotes the real part of  $z$ , that is,  $x$  when  $z = x + iy$  with  $x$  and  $y$  real.) Must choose  $c$  so that  $M(c) < \infty$ . Rewrite inversion formula using cumulant generating function  $K(t) = \log(M(t))$ :

$$2\pi i f(x) = \int_{c-i\infty}^{c+i\infty} \exp(K(z) - zx) dz.$$

Along the contour in question we have  $z = c + iy$  so we can think of the integral as being

$$i \int_{-\infty}^{\infty} \exp(K(c + iy) - (c + iy)x) dy.$$

Now do a Taylor expansion of the exponent:

$$\begin{aligned} K(c + iy) - (c + iy)x = \\ K(c) - cx + iy(K'(c) - x) - y^2 K''(c)/2 + \dots \end{aligned}$$

Ignore the higher order terms and select a  $c$  so that the first derivative

$$K'(c) - x$$

vanishes. Such a  $c$  is a saddlepoint. We get the formula

$$\begin{aligned} 2\pi f(x) \approx \exp(K(c) - cx) \\ \times \int_{-\infty}^{\infty} \exp(-y^2 K''(c)/2) dy. \end{aligned}$$

Integral is normal density calculation; gives

$$\sqrt{2\pi/K''(c)}.$$

Saddlepoint approximation is

$$f(x) = \frac{\exp(K(c) - cx)}{\sqrt{2\pi K''(c)}}.$$

Essentially same idea: Laplace's approximation.

**Example:** Sterling's approximation to factorial:

$$n! = \int_0^{\infty} \exp(n \log(x) - x) dx .$$

Exponent maximized when  $x = n$ .

For  $n$  large approximate  $f(x) = n \log(x) - x$  by  $f(x) \approx f(x_0) + (x - x_0)f'(x_0) + (x - x_0)^2 f''(x_0)/2$  and choose  $x_0 = n$  to make  $f'(x_0) = 0$ . Then

$$n! \approx \int_0^{\infty} \exp[n \log(n) - n - (x - n)^2 / (2n)] dx .$$

Substitute  $y = (x - n)/\sqrt{n}$ ; get approximation

$$n! \approx n^{1/2} n^n e^{-n} \int_{-\infty}^{\infty} e^{-y^2/2} dy$$

or

$$n! \approx \sqrt{2\pi n} n^{n+1/2} e^{-n} .$$

Note: sloppy about limits of integration.

Real proof must show integral over  $x$  not near  $n$  is negligible.